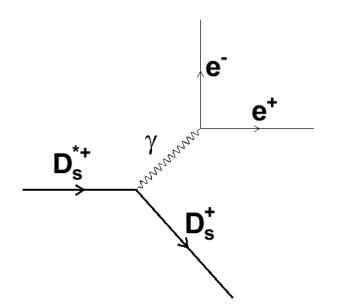


Cornell University Laboratory for Elementary-Particle Physics



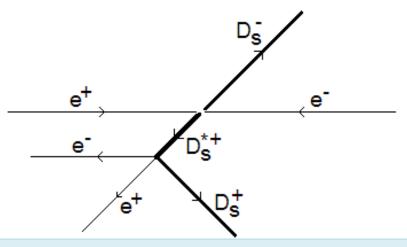
OBSERVATION OF THE DECAY $D_s^{*+} \rightarrow D_s^+ e^+ e^-$ AND MEASUREMENT OF THE RATIO OF BRANCHING FRACTIONS $B(D_s^{*+} \rightarrow D_s^+ e^+ e^-)/B(D_s^{*+} \rightarrow D_s^+ \gamma)$ AT THE CLEO-C EXPERIMENT

> Souvik Das Cornell University



10 December 2010

What Are We Looking For?



•Searching for $D_s^{*+} \rightarrow D_s^{+} e^+ e^-$ with a blind analysis.

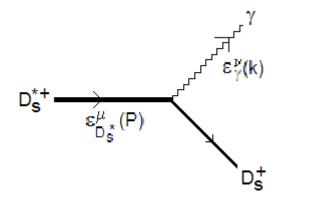
•Known decay channels are:

• $D_s^{*+} \rightarrow D_s^{+} \gamma$; Branching Fraction = 94.2% • $D_s^{*+} \rightarrow D_s^{+} \pi^0$; Branching Fraction = 5.8% [1]

•We are using e^+e^- collision data collected by the CLEO-c detector at the Cornell Electron Storage Ring (CESR) operating at $\sqrt{s} = 4170$ MeV. We have $586 \pm 6 \text{ pb}^{-1}$ of data at this energy.

• $D_s^{*\pm} D_s^{\mp}$ Production cross section at this energy is 948 ± 36 pb (combining results from [2] and [3]). This will give us ~ 555,000 events to work with.

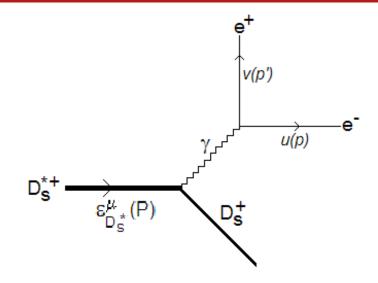
Predicted $D_s^{*\pm} \rightarrow D_s^{\pm} e^+ e^-$ Rate



If we write the matrix element of the D_s^{*+} decay to a real photon in the form:

 $M = \varepsilon^{\mu}_{D^{*+}_{S}} \varepsilon^{*\nu}_{\gamma} T_{\mu\nu}(P,k)$

where $T_{\mu\nu}(P,k)$ is a generic form factor coupling the D_s^{*+} with a photon.



Then we can write the matrix element of the decay to e^+e^- in the form:

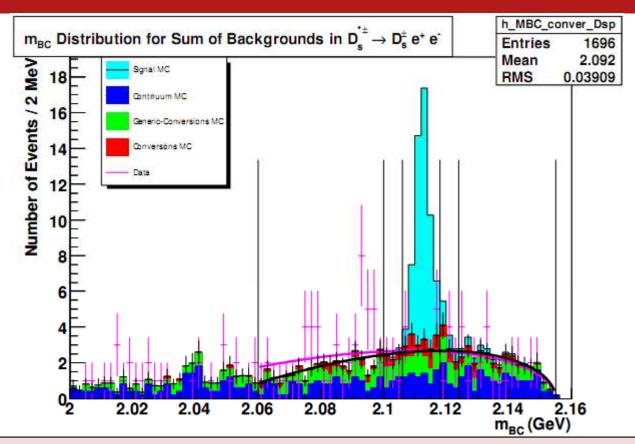
$$M = \varepsilon_{D_{S}^{*+}}^{\mu} T_{\mu\nu}(P,k) \left(\frac{-ig^{\nu\sigma}}{k^{2}}\right) \overline{u}(p) ie\gamma_{\sigma} v(p')$$

Evaluating the spin-average over the initial states and spin-sum over the final states of the invariant amplitudes and integrating over the phase space of daughters, we predict the ratio of branching fractions:

$$\frac{B(D_s^{*+} \rightarrow D_s^+ e^+ e^-)}{B(D_s^{*+} \rightarrow D_s^+ \gamma)} = 0.65\% = 0.89\alpha$$

3

Estimation of Background from m_{BC} Sidebands



•We try to be as Monte Carlo independent as possible. Extrapolate data from sidebands into signal region. Trying to fit a shape to the individual modes is impossible due to low statistics.

•We add all modes and plot the m_{BC} as shown.

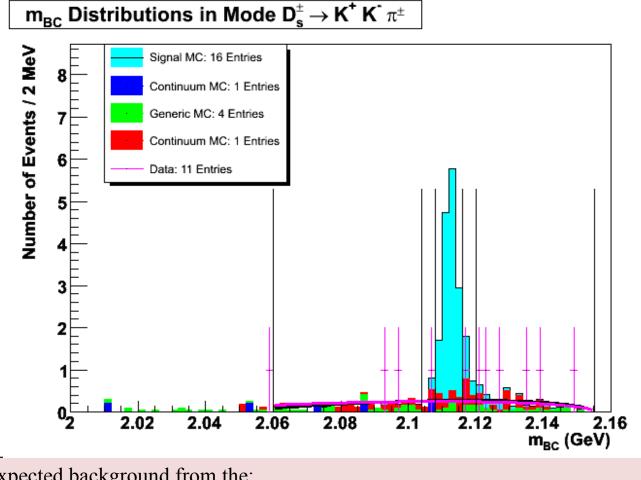
•Determine a fit shape from

•Monte Carlo backgrounds (black curve)

•Data in the sidebands (signal region is still blind!) (pink curve)

•Fit these curves in the individual channels to estimate background in the signal region.

Estimation of Background from m_{BC} Sidebands in the $K^+K^-\pi^+$ Mode



•Expected background from the:

•Monte Carlo shape (black curve): 1.10 ± 0.39 (stat) events

•Data shape (pink curve): 1.0 ± 0.35 (stat) events

•Average: 1.05 ± 0.37 (stat) events. Quoted as estimated background.

•Repeat procedure with δm distributions to obtain systematic uncertainty:

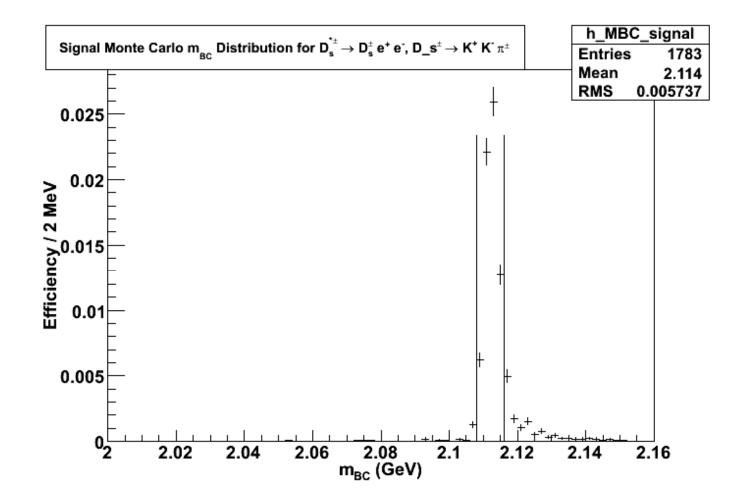
 -1.05 ± 0.37 (stat) ± 0.79 (syst)

Prediction for Signal from Monte Carlo and Data Sidebands

Decay Mode of the <i>D</i> _{<i>S</i>} ⁺	Expected Yield in 586 pb ⁻¹ (Signal + Background)	Expected Background in 586 pb ⁻¹
$K^+K^-\pi^+$	14.70	$1.05 \pm 0.37 \pm 0.79$
$K_s K^+$	3.87	$0.85 \pm 0.43 \pm 0.74$
$\pi^+\eta; \eta { ightarrow} \gamma\gamma$	3.21	$1.40 \pm 0.70 \pm 0.49$
$\pi^+\dot{\eta};\dot{\eta}{ ightarrow}\pi^+\pi^-\eta;\eta{ ightarrow}\gamma\gamma$	1.20	0.00 + 0.63 + 0.00
$K^{\scriptscriptstyle +}K^{\scriptscriptstyle -}\pi^{\scriptscriptstyle +}\pi^0$	6.55	$1.70 \pm 0.47 \pm 0.56$
$\pi^+\pi^-\pi^+$	5.32	$1.57 \pm 0.45 \pm 0.59$
$K^{*+}K^{*0}$; $K^{*+} \rightarrow K^0{}_S\pi^+$; $K^{*0} \rightarrow K^-\pi^+$	3.57	$1.58 \pm 0.53 \pm 0.40$
$\eta ho^+;\eta{ ightarrow}\gamma\gamma; ho^+{ ightarrow}\pi^+\pi^0$	8.11	$2.62 \pm 0.59 \pm 0.23$
$\acute{\eta}\pi^+$; $\acute{\eta}{ ightarrow} ho^0$ γ	4.26	$1.84 \pm 0.49 \pm 0.25$
Total	50.79	$12.61 \pm 1.58 \pm 1.53$

If $D_s^{*+} \rightarrow D_s^{+}e^+e^-$ exists, and our QED based estimation of its rate is correct, we should see a clear signal over the background for it in our data on unblinding.

Signal Selection Efficiencies



•What fraction of produced signal events are retained by our selection criteria? This was calculated from the signal region in the m_{BC} distribution.

Calculating the Ratio of Branching Fractions

For the *i*th decay mode of the D_s^+ , we may write an expression for the signal yield N_{e+e}^i .

$$L\sigma B(D_s^{*+} \to D_s^+ e^+ e^-) B(D_s^+ \to i) \varepsilon_{D_s e^+ e^-}^i = N_{e^+ e^-}^i$$

Where *L* is the integrated luminosity of data used (586 pb⁻¹), σ is the production cross section of $D_s^{*+}D_s^{-}$ at 4170 MeV (948 pb), and ε_{Dse+e}^{i} is the selection efficiency of criteria for the *i*th mode.

We can write a similar expression for the $D_s^{*+} \rightarrow D_s^{+} \gamma$ channel:

$$L\sigma B(D_s^{*+} \to D_s^{+}\gamma)B(D_s^{+} \to i)\varepsilon_{D_s\gamma}^{i} = N_{\gamma}^{i}$$

We will evaluate the ratio of branching fractions with all modes considered together:

$$K = \frac{B(D_s^{*+} \to D_s^+ e^+ e^-)}{B(D_s^{*+} \to D_s^+ \gamma)} = \left(\frac{\sum_i N_{e+e-}^i}{\sum_i N_{\gamma}^i}\right) \left(\frac{\sum_i \varepsilon_{D_s \gamma}^i B(D_s^+ \to i)}{\sum_i \varepsilon_{D_s e+e-}^i B(D_s^+ \to i)}\right)$$

Signal Selection Efficiencies

To see the contribution of uncertainties to this measure of K, we may write it as

$$K = \left(\frac{\sum_{i} N_{e+e^{-}}^{i}}{\sum_{i} N_{\gamma}^{i}}\right) \left(\frac{\varepsilon_{\gamma}}{\varepsilon_{e+e^{-}}}\right) \left(\frac{\sum_{i} \varepsilon_{D_{s}}^{i} B(D_{s}^{+} \to i)}{\sum_{i} \varepsilon_{D_{s}}^{i} B(D_{s}^{+} \to i)}\right) \Longrightarrow K = \left(\frac{\sum_{i} N_{e+e^{-}}^{i}}{\sum_{i} N_{\gamma}^{i}}\right) \left(\frac{\varepsilon_{\gamma}}{\varepsilon_{e+e^{-}}}\right)$$

To be fair, the ε_{Ds}^{i} in the numerator and denominator are not exactly equal because slightly different selection criteria are used for reconstructing $D_{s}^{*+} \rightarrow D_{s}^{+}e^{+}e^{-}$ and $D_{s}^{*+} \rightarrow D_{s}^{+}\gamma$ but they are close enough to assume the last term to be 1.

The uncertainties in *K* may be broken down as follows.

$$\left(\frac{\Delta K(stat)}{K}\right)^{2} = \frac{\sum_{i} \left(\Delta N_{e+e-}^{i}(stat)\right)^{2}}{\left(\sum_{i} N_{e+e-}^{i}\right)^{2}} + \frac{\sum_{i} \left(\Delta N_{\gamma}^{i}(stat)\right)^{2}}{\left(\sum_{i} N_{\gamma}^{i}\right)^{2}}$$
$$\left(\frac{\Delta K(syst)}{K}\right)^{2} = \frac{\sum_{i} \left(\Delta N_{e+e-}^{i}(syst)\right)^{2}}{\left(\sum_{i} N_{e+e-}^{i}\right)^{2}} + \frac{\sum_{i} \left(\Delta N_{\gamma}^{i}(syst)\right)^{2}}{\left(\sum_{i} N_{\gamma}^{i}\right)^{2}} + \left(\frac{\Delta \left(\varepsilon_{\gamma} / \varepsilon_{e^{+}e^{-}}\right)}{\varepsilon_{\gamma} / \varepsilon_{e^{+}e^{-}}}\right)^{2}}$$

Criteria to Select $D_s^{*+} \rightarrow D_s^{+} \gamma$ Events

•We also measure the signal yields and efficiencies for the $D_s^{*+} \rightarrow D_s^{+} \gamma$ channel since we have set out to measure the ratio of branching fractions

$$B(D_s^{*+} \to D_s^+ e^+ e^-) / B(D_s^{*+} \to D_s^+ \gamma)$$

•We reconstruct the D_s^{*+} through the D_s^{+} and the γ

•The D_s^+ is reconstructed through the 9 hadronic decay modes

•Criteria on the e^+e^- inapplicable. Instead, criteria on the electromagnetic shower of the photon in the calorimeter are applied.

•10 MeV < Shower Energy < 2.0 GeV

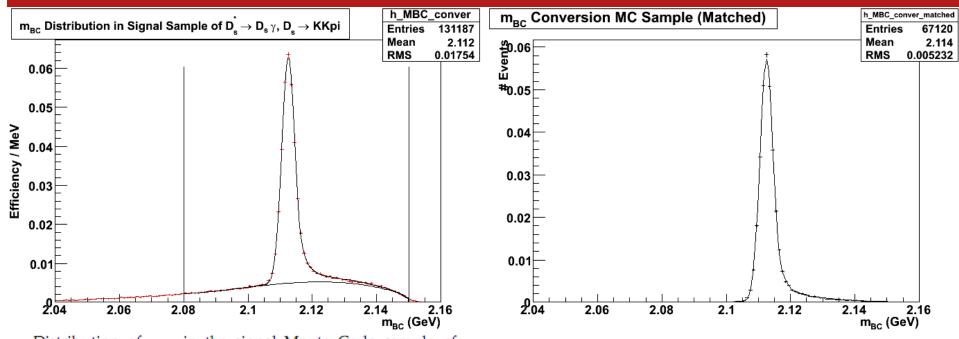
•No tracks leading to or in the vicinity of the shower

•Known noisy calorimeter crystals discarded

•Electromagnetic shower shape

•The m_{BC} criterion is discarded in favor of a fit to extract yield.

Criteria to Select $D_s^{*+} \rightarrow D_s^{+} \gamma$ Events where $D_s^{+} \rightarrow K^+ K^- \pi^+$



Distribution of m_{BC} in the signal Monte Carlo sample of $D_s^{*+} \rightarrow D_s^+ \gamma$ events where $D_s^+ \rightarrow K^+ K^- \pi^+$. The plot is normalized so as to directly read out the efficiency of the m_{BC} selection criterion.

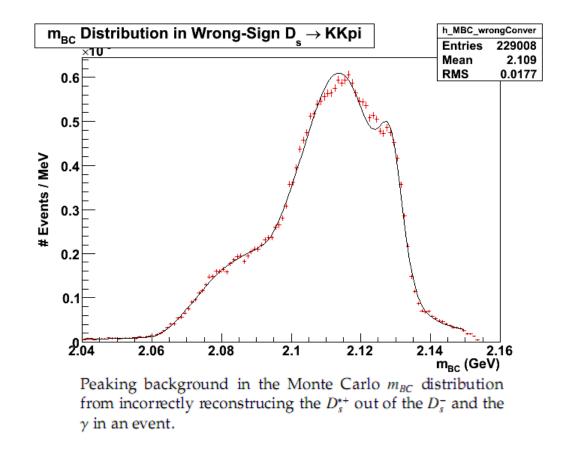
Distribution of m_{BC} in signal MC where the D_s^+ and the γ have been MC matched.

•Shape of the peak determined from events where both the Ds+ and photon are MC matched. Used a Crystal Ball function with power law on high side with soft Gaussian on high shoulder.

•Fitted peak on top of: $f(x; x_0, p, C_0, C_1, C_2, C_3) = (C_0 + C_1 x + C_2 x^2 + C_3 x^3)(x - x_0)^p; 0$

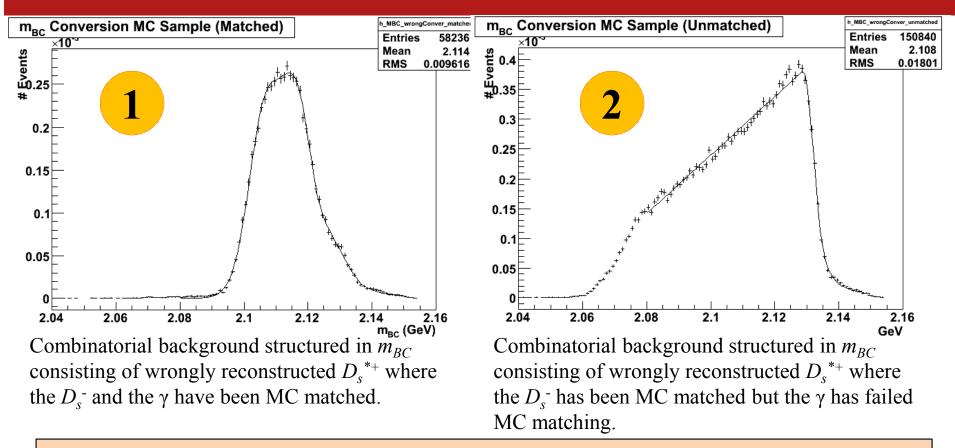
•Extracted signal efficiency

Criteria to Select $D_s^{*+} \rightarrow D_s^+ \gamma$ Events where $D_s^+ \rightarrow K^+ K^- \pi^+$



•Incorrectly reconstructed D_s^{*+} mesons from the combination of D_s^{-} and γ exhibit a double-humped structure in the m_{BC} distribution.

Criteria to Select $D_s^{*+} \rightarrow D_s^{+} \gamma$ Events where $D_s^{+} \rightarrow K^+ K^- \pi^+$

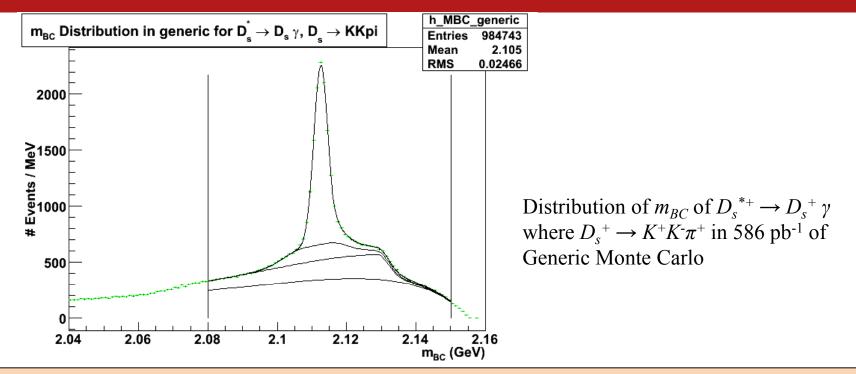


•Shape 1 was fitted to a Gaussian on the left, a Crystal Ball function on the right with its power law on the high side and another soft Gaussian on the high shoulder.

•Shape 2 was fitted to a Crystal Ball function with its power law on the high side and analytically continued into a straight line on the low side.

•The amplitudes of these shapes were fitted independently to data and generic MC.

Criteria to Select $D_s^{*+} \rightarrow D_s^{+} \gamma$ Events where $D_s^{+} \rightarrow K^+ K^- \pi^+$



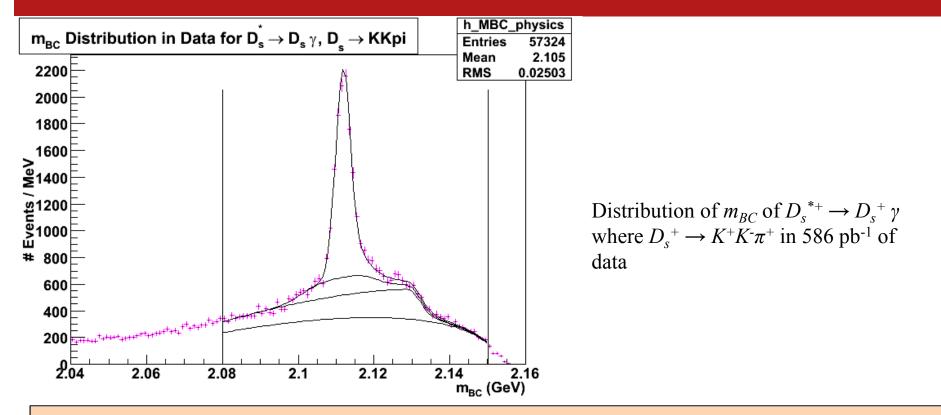
•Try fit Generic MC to see if we can recover $B(D_s^{*+} \rightarrow D_s^{+} \gamma)$

•Fits from bottom to top:

- 1. Featureless combinatorial background modeled by an Argus function.
- 2. Scaled shape 2, wrongly reconstructed D_s^{*+} using a matched D_s^{-} and a matched photon
- 3. Scaled shape 1, wrongly reconstructed D_s^{*+} using a matched D_s^{-} and a random photon
- 4. Scaled peak shape recovered from signal MC where D_s^+ and the photon were matched

• $B(D_s^{*+} \rightarrow D_s^{+} \gamma) = 0.926 \pm 0.006$. [Systematics arising from inconsistencies between models used for signal and generic MC not estimated.]

Criteria to Select $D_s^{*+} \rightarrow D_s^+ \gamma$ Events where $D_s^+ \rightarrow K^+ K^- \pi^+$



•Signal selection efficiency = 0.339 ± 0.002

•Signal yield = $9114 \pm 110 \pm 201$

• $B(D_s^{*+} \rightarrow D_s^{+} \gamma) = 0.880 \pm 0.011^{[1]} \pm 0.045^{[2]} \pm 0.035^{[3]} \pm 0.019^{[4]}$

•[1] is the statistical uncertainty from the fit

•[2] is the systematic uncertainty from the uncertainty in $B(D_s^+ \to K^+ K^- \pi^+)$.

•[3] encapsulates the systematic uncertainty from the signal efficiency, the integrated luminosity and production cross section of $D_s^{*+}D_s^{-}$

•[4] is the systematic uncertainty from the fit. Evaluated using an alternative fitting method. •Our measurement of $B(D_s^{*+} \rightarrow D_s^{+} \gamma)$ is 1σ away from the accepted value of 0.942 ± 0.007

Yields, Efficiencies and Unblinding Data

•The yields and efficiencies for $D_s^{*+} \rightarrow D_s^{+} \gamma$ in the 9 hadronic decay modes of the Ds+ are recorded. Will be presented in final calculation of

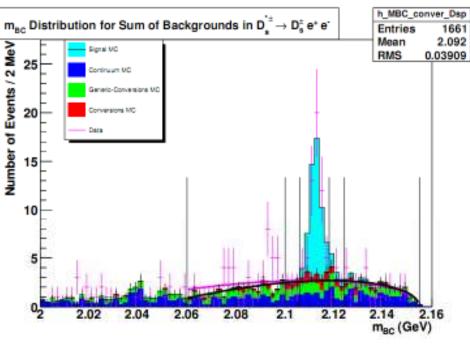
$$\frac{B(D_s^{*+} \to D_s^+ e^+ e^-)}{B(D_s^{*+} \to D_s^+ \gamma)}$$

where systematic uncertainties in the reconstruction of Ds+ mesons would have canceled.

But now we unblind the signal region for $D_s^{*+} \rightarrow D_s^+ e^+ e^-$ So buckle your seat belts!

Unblinding Data in the m_{BC} Signal Region of $D_s^{*+} \rightarrow D_s^{+} e^+ e^-$

Decay Mode of the D_S^+	Expected Yield from MC in 586 pb ⁻¹	Observed Yield in 586 pb ⁻¹
$K^+K^-\pi^+$	14.70	14
$K_s K^+$	3.87	1
$\pi^+\eta; \eta { ightarrow} \gamma\gamma$	3.21	4
$\pi^+\dot{\eta};\dot{\eta}{ ightarrow}\pi^+\pi^-\eta;\eta{ ightarrow}\gamma\gamma$	1.20	4
$K^{\scriptscriptstyle +}K^{\scriptscriptstyle -}\pi^{\scriptscriptstyle +}\pi^0$	6.55	6
$\pi^+\pi^-\pi^+$	5.32	7
$K^{*+}K^{*0}; K^{*+} \longrightarrow K^0{}_S\pi^+; K^{*0} \longrightarrow K^-\pi^+$	3.57	4
$\eta ho^+;\eta{ o}\gamma\gamma; ho^+{ o}\pi^+\pi^0$	8.11	7
$\acute{\eta}\pi^+$; $\acute{\eta}{ ightarrow} ho^0\gamma$	4.26	4
Total	50.79	51



•Distribution of m_{BC} in data after unblinding all modes.

•Magenta curve extrapolates to estimate background using the data shape. Black curve extrapolates data using the MC shape.

Results from Unblinding Data

Decay Mode of the <i>D</i> _S ⁺	Branching Fraction	$D_s^{*+} \rightarrow D_s^+ e^+ e^-$		$D_s^{*+} \rightarrow D_s^+ \gamma$		$\mathcal{D}(\mathcal{D}^{*+} \rightarrow \mathcal{D}^{+} - 0)$
(RPP 2008) %	Background Subtracted Yield	Selection Efficiency	Background Subtracted Yield	Selection Efficiency	$K = \frac{B(D_s^{*+} \to D_s^{+} \pi^0)}{B(D_s^{*+} \to D_s^{+} \gamma)}$	
$K^+K^-\pi^+$	0.055 ± 0.0028	$12.95 \pm 3.76 \pm 0.79$	0.0730 ± 0.0019	$9114\pm110\pm201$	0.339 ± 0.002	$0.0066 \pm 0.0019 \pm 0.0005$
$K_s K^+$	0.0149 ± 0.0009	$0.15 \pm 1.09 \pm 0.74$	0.0597 ± 0.0017	$1902\pm57\pm45$	0.2573 ± 0.0004	$0.0003 \pm 0.0025 \pm 0.0017$
$\pi^+\eta;\ \eta{ ightarrow}\gamma\gamma$	0.0062 ± 0.0008	$2.60 \pm 2.12 \pm 0.49$	0.0855 ± 0.0021	$1037\pm46\pm\!\!37$	0.3310 ± 0.0015	$0.0097 \pm 0.0079 \pm 0.0019$
$\pi^+ \eta; \ \eta \longrightarrow \pi^+ \pi^- \eta; \ \eta \longrightarrow \gamma \gamma$	0.0067 ± 0.0007	$4.00 \pm 2.10 \pm 0.00$	0.0530 ± 0.0016	$691 \pm 34 \pm 40$	0.2101 ± 0.0013	$0.023 \pm 0.0123 \pm 0.0015$
$K^+K^-\pi^+\pi^0$	0.056 ± 0.005	$4.30 \pm 2.49 \pm 0.56$	0.0255 ± 0.0011	$3592\pm118\pm72$	0.1225 ± 0.0010	$0.0058 \pm 0.0033 \pm 0.0008$
$\pi^+\pi^-\pi^+$	0.0111 ± 0.0008	$5.43 \pm 2.68 \pm 0.59$	0.0992 ± 0.0022	$2745\pm93\pm52$	0.4583 ± 0.0018	$0.0091 \pm 0.0045 \pm 0.0010$
$K^{*+}K^{*0};$ $K^{*+} \rightarrow K^{0}{}_{S}\pi^{+};$ $K^{*0} \rightarrow K^{-}\pi^{+}$	0.0164 ± 0.0012	$2.42 \pm 2.07 \pm 0.40$	0.0356 ± 0.0013	$1570 \pm 74 \pm 13$	0.1913 ± 0.0012	$0.0083 \pm 0.0071 \pm 0.0014$
$\begin{array}{c} \eta ho^+;\ \eta o \gamma\gamma;\ ho^+ o \pi^+\pi^0 \end{array}$	0.0511 ± 0.0087	$4.38 \pm 2.71 \pm 0.23$	0.0316 ± 0.0013	$3170 \pm 161 \pm 313$	0.1839 ± 0.0013	$0.0080 \pm 0.0050 \pm 0.0010$
$\acute{\eta}\pi^+$; $\acute{\eta} ightarrow ho^0$ γ	0.0112 ± 0.0012	$2.16 \pm 2.06 \pm 0.25$	0.064 ± 0.0018	$1531\pm80\pm122$	0.3171 ± 0.0015	$0.0070 \pm 0.0067 \pm 0.0010$
Sum of Modes		$38.39 \pm 7.32 \pm 1.53$		25351 ± 280		$0.0072 \pm 0.0014 \pm 0.0003$

Result

•We observe a signal for $D_s^{*+} \rightarrow D_s^+ e^+ e^-$ with a signal significance of 6.39 σ

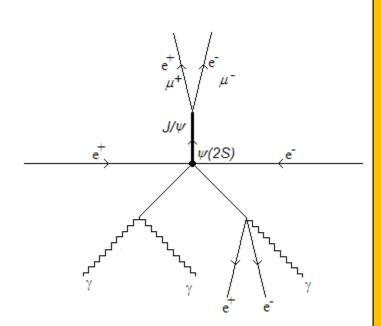
•We measure the ratio of branching fractions from the table:

$$K = \frac{B(D_s^{*+} \to D_s^+ e^+ e^-)}{B(D_s^{*+} \to D_s^+ \gamma)} = (0.72 \pm 0.14(stat) + 0.03(syst))\%$$

•However, the multiplicative systematic from the last term of $\Delta K(syst)$ on Slide 9 has not been included. The fractional error from this term is measured by an independent study to be 6.51%. 6.51% of 0.72% is 0.047% and when added in quadrature to the other systematics, we arrive at the final result

 $\frac{B(D_s^{*+} \to D_s^+ e^+ e^-)}{B(D_s^{*+} \to D_s^+ \gamma)} = (0.72 \pm 0.14(stat) + 0.06(syst))\%$

•The systematic uncertainty stemming from uncertainties in tracking soft electrons and photons will contribute multiplicatively:



 $\Delta K = K \frac{\Delta \left(\varepsilon_{\gamma} / \varepsilon_{e^+ e^-} \right)}{\varepsilon_{\gamma} / \varepsilon_{e^+ e^-}}$

•We estimate this by measuring $B(\pi^0 \rightarrow e^+e^-\gamma)$ within the energy range of our analysis. Discrepancies between this measurement and the currently accepted value is related to the systematic uncertainty we are trying to establish.

 $\frac{\Delta \left(\varepsilon_{\gamma} / \varepsilon_{e^+ e^-} \right)}{\varepsilon_{\gamma} / \varepsilon_{e^+ e^-}} = \frac{\Delta B(\pi^0 \to e^+ e^- \gamma)}{B(\pi^0 \to e^+ e^- \gamma)}$

•To study these π^0 , we look at $\psi(2S) \rightarrow J/\psi \pi^0 \pi^0$ events.

- •We call events where one of the π^0 Dalitz decays to $\pi^0 \rightarrow \gamma e^+ e^- Type I$ events
- Events where both $\pi^0 \rightarrow \gamma \gamma$ (*Type II* events)

$$\frac{n_I}{n_{II}} = \frac{2B(\pi^0 \to e^+ e^- \gamma)}{B(\pi^0 \to \gamma \gamma)}$$
²⁰

•We employ two methods to estimate $B(\pi^0 \rightarrow e^+e^-\gamma)$ in order to assign a systematic uncertainty to our measurement.

METHOD 1

We set up selection criteria to reconstruct the ψ(2S) from Type I events.
Selection efficiency to keep Type I events from an MC sample is called ε_s. Efficiency to keep Type II events from an MC sample is called ε_c.
The yield in data, y, is related to the numbers of produced Type I and II by

 $n_I \varepsilon_s + n_{II} \varepsilon_c = y$

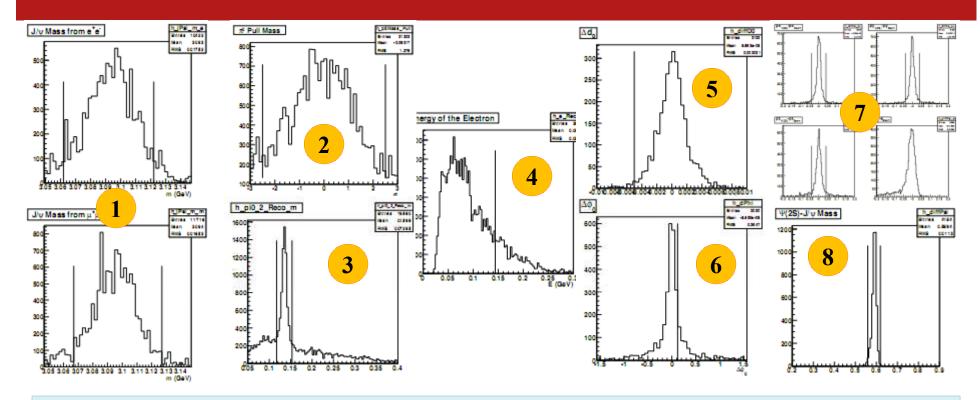
•We solve for n_I keeping in mind the currently accepted ratio of n_I and n_{II}

We set up selection criteria to reconstruct the ψ(2S) from Type I events.
Applying them to our data, we find their signal yield and directly estimate the number of Type II events in our data.

•Having found n_I and n_{II} , we calculate $B(\pi^0 \rightarrow \gamma e^+ e^-)$ from

$$\frac{n_I}{n_{II}} = \frac{2B(\pi^0 \to e^+ e^- \gamma)}{B(\pi^0 \to \gamma \gamma)}$$

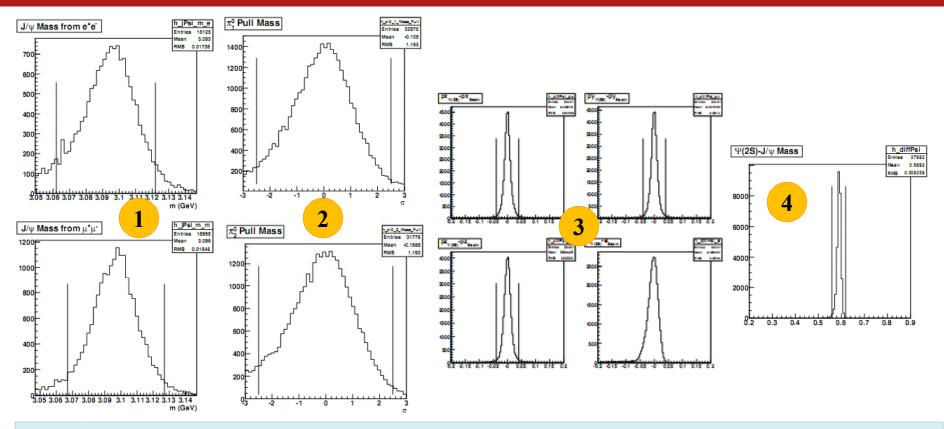
Selection Criteria for Type I Events



- The J/ψ is reconstructed from e^+e^- or $\mu^+\mu^-$ 1.
- First π^0 reconstructed from two photons. Criterion on pull mass 2.
- Second π^0 reconstructed from one photon and a soft e^+e^- pair 3.
- The energy of the e+e- are restricted to within 144 MeV. 4.
- Δd_0 cuts applied on the soft e+e- pair 5.
- $\Delta \varphi_0$ cuts applied on the soft e+e- pair 6.
- The $\psi(2S)$ is required to be within a range of the collision lab 4-momentum 7.

8. $m_{\psi(2S)} - m_{J/\psi}$ is required to be within a range. •We find ε_s and ε_c using these on MC and apply them on data to record the yield, *y*.

Selection Criteria for Type II Events



- 1. The J/ψ is reconstructed from e^+e^- or $\mu^+\mu^-$
- 2. Both π^0 reconstructed from two photons. Criterion on pull mass
- 3. The $\psi(2S)$ is required to be within a range of the collision lab 4-momentum
- 4. $m_{\psi(2S)}$ $m_{J/\psi}$ is required to be within a range.

•We find the efficiency of keeping Type II events using these.

•We record the yield in data.

METHOD 1

•Solving for n_I and n_{II} , we get

$$\frac{n_I}{n_{II}} = \frac{8447 \pm 554}{341607 \pm 2555} = \frac{2B(\pi^0 \to e^+ e^- \gamma)}{(98.823 \pm 0.034) \times 10^{-2}}$$

•This gives us $B(\pi^0 \rightarrow e^+ e^- \gamma) = 0.01222 \pm 0.00081$ (stat)

METHOD 2

•We set up selection criteria to reconstruct the $\psi(2S)$ from Type II events where the photon likely converted to an e^+e^- . This is done by accepting events which were formerly rejected by the Δd_0 or $\Delta \varphi_0$ selection criteria.

•Selection efficiency to keep Type I events from an MC sample is called ε'_{s} . Efficiency to keep Type II events from an MC sample is called ε'_{c} . The yield in data, y', is related to the numbers of produced Type I and II by

$$n_I \varepsilon'_s + n_{II} \varepsilon'_c = y'$$

•Using this equation simultaneously with that on Slide 21, we solve for n_I . We use numbers for n_I as derived before.

$$\frac{n_I}{n_{II}} = \frac{8437 \pm 342}{341607 \pm 2555} = \frac{2B(\pi^0 \to e^+ e^- \gamma)}{(98.823 \pm 0.034) \times 10^{-2}}$$

•From this we get $B(\pi^0 \rightarrow e^+e^-\gamma) = 0.01220 \pm 0.00050$ (stat)

•Combining Method 1 & 2, $B(\pi^0 \rightarrow e^+e^-\gamma) = 0.01222 \pm 0.00081 \text{ (stat)} \pm 0.00002 \text{ (syst)}$

•Currently accepted value $B(\pi^0 \rightarrow e^+e^-\gamma) = 0.01174 \pm 0.00035$

•Difference between our result and current value is 0.00046. Of the same order of magnitude as the uncertainties, and therefore added in quadrature to get a total uncertainty of 0.00077.

•Now we can write

$$\frac{\Delta\left(\varepsilon_{\gamma}/\varepsilon_{e^{+}e^{-}}\right)}{\varepsilon_{\gamma}/\varepsilon_{e^{+}e^{-}}} = \frac{\Delta B(\pi^{0} \to e^{+}e^{-}\gamma)}{B(\pi^{0} \to e^{+}e^{-}\gamma)} = \frac{0.00077}{0.01174} = 0.0651$$

and this is what we propagate into the systematic error for K

Bibliography

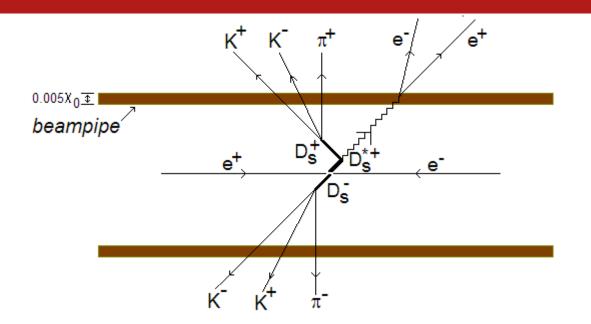
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- Werner Sun for helping me familiarize myself with CLEO-c software, reprocess data, and offering me general physics advice when sought.
- Peter Onyisi for answering all my questions about D_s^+ mesons.
- Brian Heltsley for help with reprocessing data, answering my questions about CLEO-c calorimetry and taking so much care in reading our note and offering interesting and useful suggestions.
- Matthew Shepherd for the feedback on our note.
- Paras Naik for fielding random questions about the detector and software at equally random times of the day or night.

Backup Slides

Typical Background Processes



Photon Conversion Background

- A background that resembles the signal is expected from D_s^{*+} decaying to $D_s^{+} \gamma$ and the γ converting to e^+e^- in the beam-pipe and other material.
- Given that the beam-pipe is $\sim 0.5\%$ of a radiation length, we can estimate this conversion background to occur at roughly the same rate as the signal

Combinatorial Backgrounds

- Dalitz decay of any $\pi^0 \rightarrow \gamma e^+ e^-$ also give equally soft electrons that appear to come from interaction point
- Other combinatorial backgrounds.

Analysis Strategy

•Employ a blind analysis search for the $D_s^{*+} \rightarrow D_s^+ e^+ e^-$

•Fully reconstruct the D_s^{*+} through the D_s^{+} and e^+e^- .

•The D_s^+ is reconstructed through 9 hadronic decay modes

•Selection criteria based on the *invariant masses* of the D_s^+ and D_s^{*+} are optimized

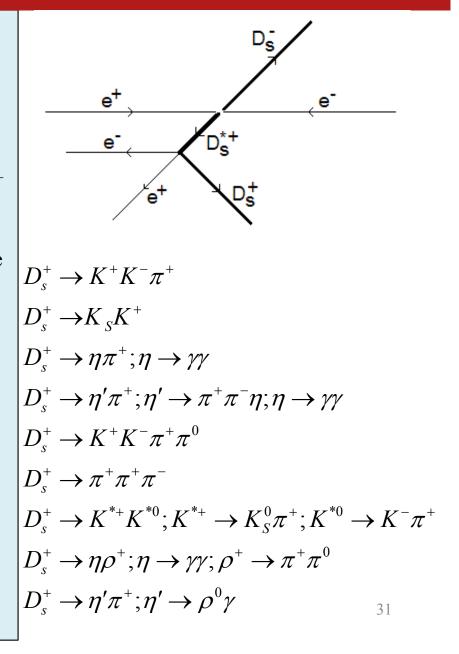
•Criteria based on the *track parameters* of the e^+ and e^- are powerful against the photon conversion background

•The e^+e^- are extremely soft (~144 MeV total) and accuracy required by tracking motives a data reprocessing campaign.

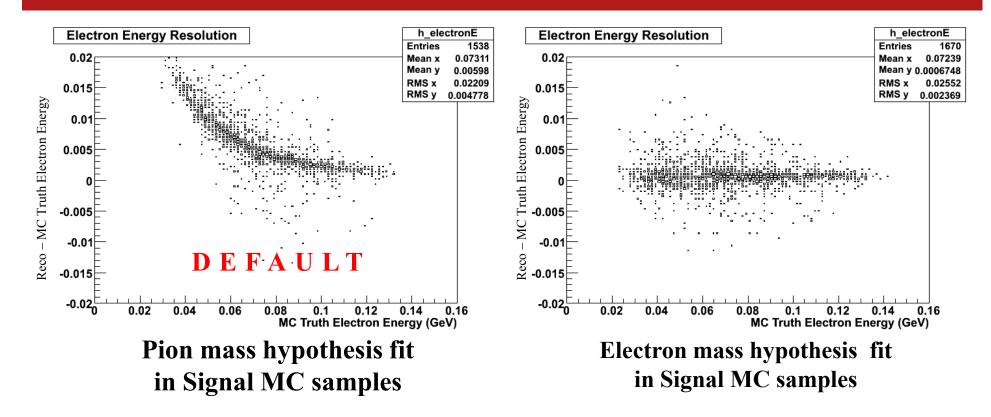
•Estimate the number of background events we will see in the signal region after our selection.

•Use selection criteria to measure the rate of $D_s^{*+} \rightarrow D_s^{+} \gamma$

•Unblind data in the signal region for $D_s^{*+} \rightarrow D_s^+ e^+ e^-$. Compute $\frac{B(D_s^{*+} \rightarrow D_s^+ e^+ e^-)}{B(D_s^{*+} \rightarrow D_s^+ \gamma)}$



Tracking Soft Electrons



- Studies with privately generated Monte Carlo samples indicate significantly better performance of analysis with electron mass hypothesis fitted data.
- Motivated a campaign to reprocess datasets containing 4170 MeV collision data to include such tracks.

Electron Track Selection Criteria

•Electron tracks must pass track quality cuts:

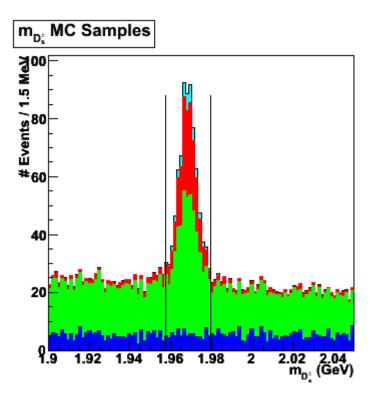
•10 MeV < Track Energy < 150 MeV

• $|z_0| < 5$ cm. Tracks pass within 5 cm of the interaction point in dimension parallel to beam axis.

• $|d_0| < 5$ mm. Tracks pass within 5 mm of the interaction point in dimensions perpendicular to the beam axis.

•dE/dx within 3.0 σ of that expected for an electron.

m_{Ds} Selection Criterion for the $K^+K^-\pi^+$ Mode



Cyan: Signal Monte Carlo Red: Conversion Monte Carlo Green: Generic Monte Carlo (cc production) Blue: Continuum Monte Carlo (light quarks)

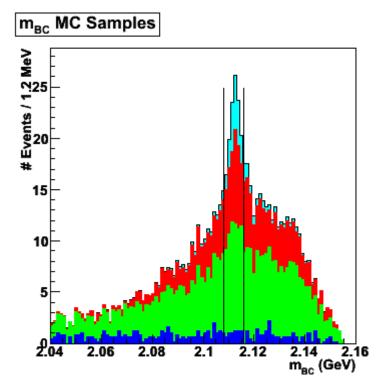
Histograms normalized to 586 pb⁻¹ of data

•We reconstruct the invariant mass m_{D_s} of a D_s^+ from its decay products.

•Selection Criterion for this mode:

$$|m_{D_s} - 1.969 GeV| < 0.011 GeV$$

m_{BC} Selection Criterion for the $K^+K^-\pi^+$ Mode



Cyan: Signal Monte Carlo Red: Conversion Monte Carlo Green: Generic Monte Carlo (cc production) Blue: Continuum Monte Carlo (light quarks)

Histograms normalized to 586 pb⁻¹

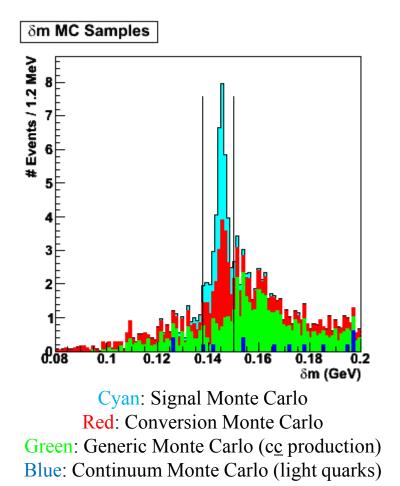
•We know the energy of the CESR beam to high precision. Given the masses of the D_s^{*+} and D_s^{+} , we can calculate the energy carried away by the D_s^{*+}

•We define the beam-constrained mass of the D_s^{*+} as:

$$m_{BC} = \sqrt{E^2 (D_S^{*+} beam) - P^2 (K^+ K^- \pi^+ e^+ e^-)}$$

•Selection Criterion for this mode: $|m_{BC} - 2.112GeV| < 0.004GeV$

δm Selection Criterion for the $K^+K^-\pi^+$ Mode



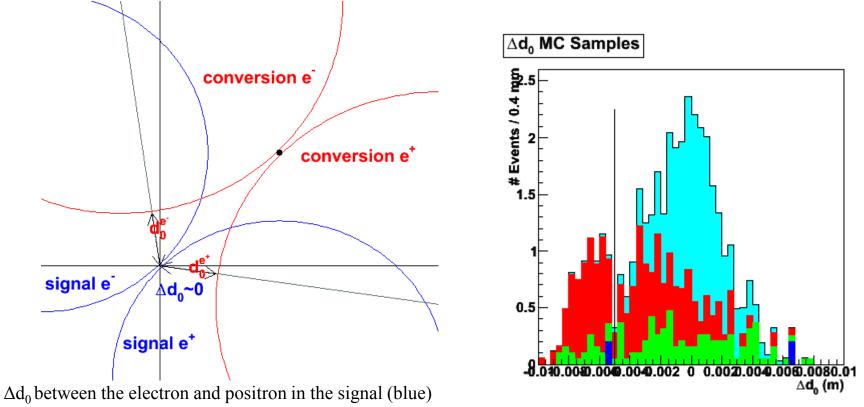
Histograms normalized to 586 pb⁻¹

•We define δm as the mass difference the D_s^{*+} and the D_s^{++} where both are reconstructed from their daughters:

 $\delta m = M(K^{+}K^{-}\pi^{+}e^{+}e^{-}) - M(K^{+}K^{-}\pi^{+})$

•Selection Criterion for this mode: $|\delta m - 0.1438 GeV| < 0.006 GeV$

Δd_0 Selection Criterion for the $K^+K^-\pi^+$ Mode



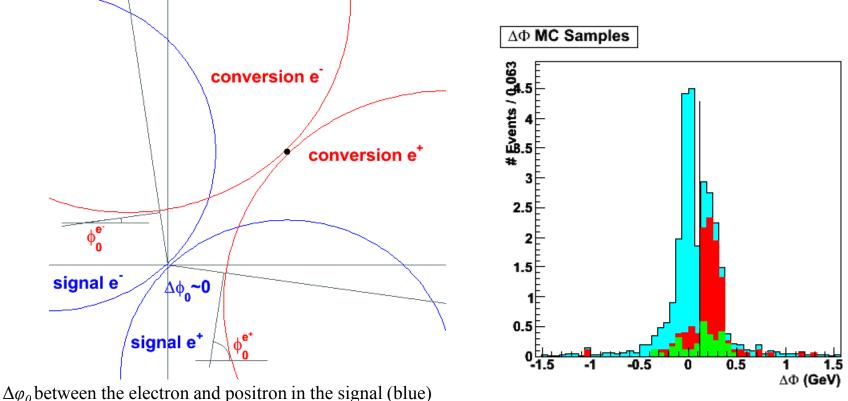
and conversion (red)

• d_{θ} : Signed distance of closest approach to the beamline [4]

•The $\Delta d_0 = d_0^{e} - d_0^{e+}$ is centered around 0 for the signal and offset from 0 for conversion backgrounds

•We require
$$\Delta d_0 > -5 \text{ mm}$$

$\Delta \varphi_0$ Selection Criterion for the $K^+ K^- \pi^+$ Mode



 $\Delta \varphi_0$ between the electron and positron in the signal (b) and conversion (red) samples

• φ_{θ} : Azimuth of track at the point of closest approach to beamline [4]

• $\Delta \varphi_0 = \varphi_0^{e} - \varphi_0^{e}$ is centered around 0 for the signal and offset for the conversion background. We require $\Delta \varphi < 0.12$

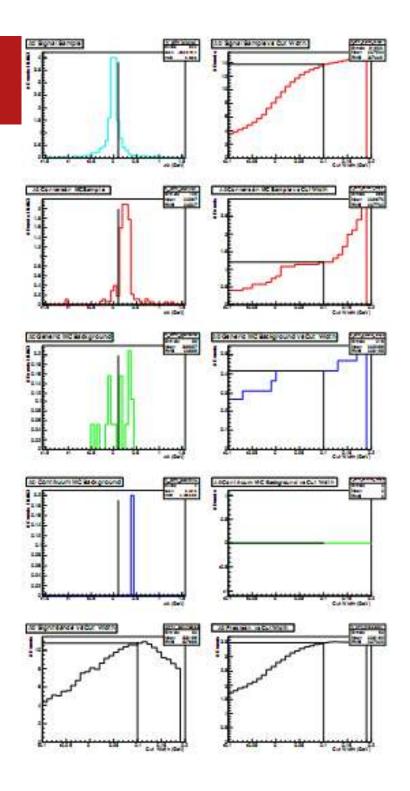
• Δd_{θ} and $\Delta \varphi_{\theta}$ constitute powerful criterion against the photon conversion background.

Optimizing Selection Criteria

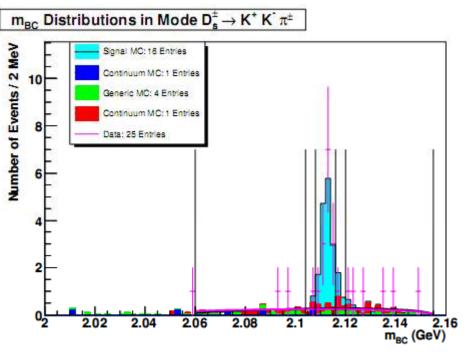
•We went channel by channel, criterion by criterion. (Example of $\Delta \Phi_0$ Selection Criterion for the $K^+K^-\pi^+$ Mode on the right)

•Plotted the signal MC, conversion MC, generic without conversion MC, and continuum MC vs variation in the cut.

•Optimized for significance $[s/\sqrt{b}]$ for low-statistics modes and precision $[s/\sqrt{(s+b)}]$ for high-statistics modes.

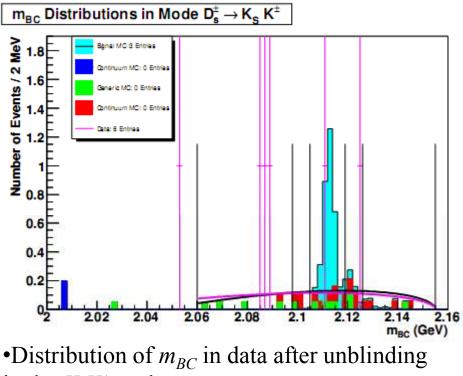


Decay Mode of the D_S^+	Expected Yield from MC in 586 pb ⁻¹	Observed Yield in 586 pb ⁻¹
$K^+K^-\pi^+$	14.70	14
$K_s K^+$	3.87	
$\pi^+\eta;\ \eta{ ightarrow}\gamma\gamma$	3.21	
$\pi^+ \dot{\eta}; \dot{\eta} { ightarrow} \pi^+ \pi^- \eta; \eta { ightarrow} \gamma \gamma$	1.20	
$K^+K^-\pi^+\pi^0$	6.55	
$\pi^+\pi^-\pi^+$	5.32	
$K^{*+}K^{*0}; K^{*+} \longrightarrow K^0{}_S\pi^+;$ $K^{*0} \longrightarrow K^-\pi^+$	3.57	
$\eta ho^+; \eta \longrightarrow \gamma \gamma; ho^+ \longrightarrow \pi^+ \pi^0$	8.11	
$\dot{\eta}\pi^+$; $\dot{\eta}{ ightarrow} ho^0$ γ	4.26	
Total	50.79	



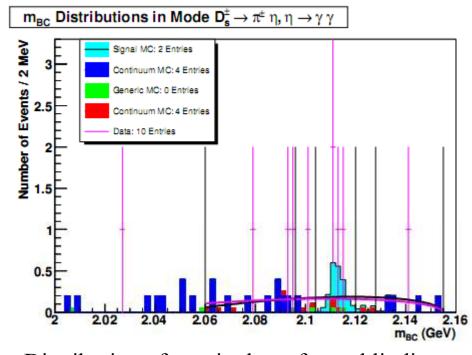
•Distribution of m_{BC} in data after unblinding in the $K^+K^-\pi^+$ mode

Decay Mode of the D_S^+	Expected Yield from MC in 586 pb ⁻¹	Observed Yield in 586 pb ⁻¹
$K^+K^-\pi^+$	14.70	14
$K_s K^+$	3.87	1
$\pi^+\eta; \eta { ightarrow} \gamma\gamma$	3.21	
$\pi^+ \dot{\eta}; \dot{\eta} {\rightarrow} \pi^+ \pi^- \eta; \eta {\rightarrow} \gamma \gamma$	1.20	
$K^{+}K^{-}\pi^{+}\pi^{0}$	6.55	
$\pi^+\pi^-\pi^+$	5.32	
$K^{*+}K^{*0}; K^{*+} \longrightarrow K^0{}_S \pi^+; K^{*0} \longrightarrow K^- \pi^+$	3.57	
$\eta ho^+; \eta ightarrow \gamma\gamma; ho^+ ightarrow \pi^+ \pi^0$	8.11	
$\acute{\eta}\pi^+$; $\acute{\eta}{ ightarrow} ho^0$ γ	4.26	
Total	50.79	



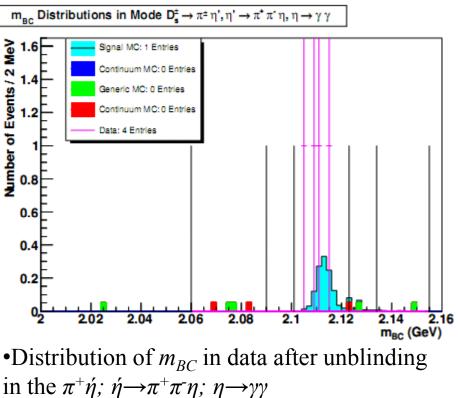
in the $K_s K^+$ mode

Decay Mode of the D_S^+	Expected Yield from MC in 586 pb ⁻¹	Observed Yield in 586 pb ⁻¹
$K^+K^-\pi^+$	14.70	14
$K_s K^+$	3.87	1
$\pi^+\eta; \eta { ightarrow} \gamma\gamma$	3.21	4
$\pi^+\dot{\eta};\dot{\eta}{ ightarrow}\pi^+\pi^-\eta;\eta{ ightarrow}\gamma\gamma$	1.20	
$K^{\scriptscriptstyle +}K^{\scriptscriptstyle -}\pi^{\scriptscriptstyle +}\pi^0$	6.55	
$\pi^+\pi^-\pi^+$	5.32	
$K^{*+}K^{*0}; K^{*+} \longrightarrow K^0{}_S\pi^+; K^{*0} \longrightarrow K^-\pi^+$	3.57	
$\eta ho^+;\ \eta o\gamma\gamma;\ ho^+ o\pi^+\pi^0$	8.11	
$\acute{\eta}\pi^+; \acute{\eta} ightarrow ho^0 \gamma$	4.26	
Total	50.79	

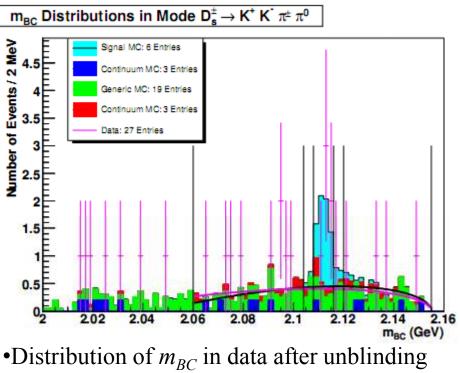


•Distribution of m_{BC} in data after unblinding in the $\pi^+\eta$; $\eta \rightarrow \gamma\gamma$ mode

Decay Mode of the D_S^+	Expected Yield from MC in 586 pb ⁻¹	Observed Yield in 586 pb ⁻¹	m _{BC} Distributions in Mo
$K^+K^-\pi^+$	14.70	14	St 1.4 Generic MC:
$K_s K^+$	3.87	1	
$\pi^+\eta; \eta { ightarrow} \gamma\gamma$	3.21	4	
$\pi^+ \dot{\eta}; \dot{\eta} { ightarrow} \pi^+ \pi^- \eta; \eta { ightarrow} \gamma \gamma$	1.20	4	0.6
$K^+K^-\pi^+\pi^0$	6.55		0.4
$\pi^+\pi^-\pi^+$	5.32		0.2
$K^{*+}K^{*0}; K^{*+} \longrightarrow K^0{}_S\pi^+;$ $K^{*0} \longrightarrow K^-\pi^+$	3.57		•Distribution of
ηho^+ ; $\eta \rightarrow \gamma \gamma$; $ ho^+ \rightarrow \pi^+ \pi^0$	8.11		in the $\pi^+ \dot{\eta}$; $\dot{\eta} \rightarrow$
$\acute{\eta}\pi^+$; $\acute{\eta} ightarrow ho^0$ γ	4.26		
Total	50.79		•Magenta curve background usi

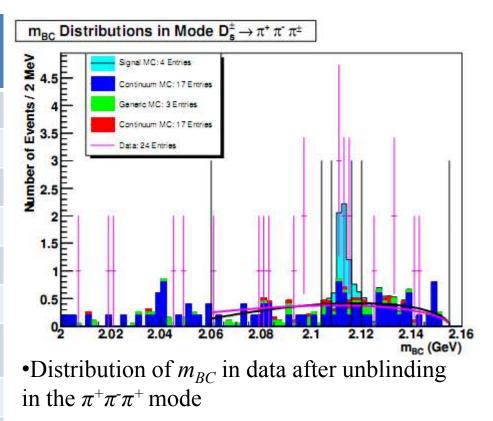


Decay Mode of the D_S^+	Expected Yield from MC in 586 pb ⁻¹	Observed Yield in 586 pb ⁻¹
$K^+K^-\pi^+$	14.70	14
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$\pi^+\eta; \eta { ightarrow} \gamma\gamma$	3.21	4
$\pi^+\dot{\eta};\dot{\eta}{ ightarrow}\pi^+\pi^-\eta;\eta{ ightarrow}\gamma\gamma$	1.20	4
$K^{\scriptscriptstyle +}K^{\scriptscriptstyle -}\pi^{\scriptscriptstyle +}\pi^0$	6.55	6
$\pi^+\pi^-\pi^+$	5.32	
$K^{*+}K^{*0}; K^{*+} \longrightarrow K^0{}_S\pi^+; K^{*0} \longrightarrow K^-\pi^+$	3.57	
$\eta ho^+;\eta{ o}\gamma\gamma; ho^+{ o}\pi^+\pi^0$	8.11	
$\acute{\eta}\pi^+; \acute{\eta} ightarrow ho^0 \gamma$	4.26	
Total	50.79	

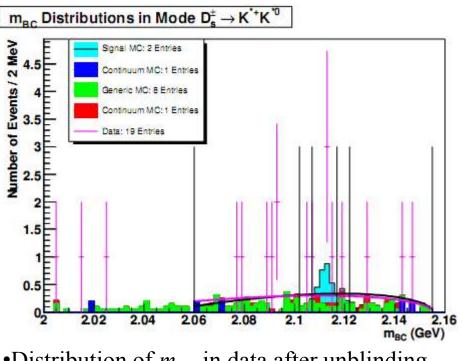


in the $K^+K^-\pi^+\pi^0$ mode

Decay Mode of the D_S^+	Expected Yield from MC in 586 pb ⁻¹	Observed Yield in 586 pb ⁻¹
$K^+K^-\pi^+$	14.70	14
$K_s K^+$	3.87	1
$\pi^+\eta; \eta { ightarrow} \gamma\gamma$	3.21	4
$\pi^+\dot{\eta};\dot{\eta}{ ightarrow}\pi^+\pi^-\eta;\eta{ ightarrow}\gamma\gamma$	1.20	4
$K^{\scriptscriptstyle +}K^{\scriptscriptstyle -}\pi^{\scriptscriptstyle +}\pi^0$	6.55	6
$\pi^+\pi^-\pi^+$	5.32	7
$K^{*+}K^{*0}; K^{*+} \longrightarrow K^0{}_S\pi^+; K^{*0} \longrightarrow K^-\pi^+$	3.57	
$\eta ho^+; \eta ightarrow \gamma\gamma; ho^+ ightarrow \pi^+ \pi^0$	8.11	
$\acute{\eta}\pi^+; \acute{\eta} ightarrow ho^0 \gamma$	4.26	
Total	50.79	

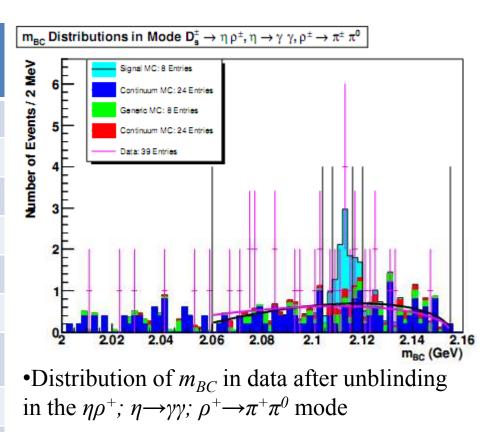


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$K^{\scriptscriptstyle +}K^{\scriptscriptstyle -}\pi^{\scriptscriptstyle +}\pi^0$	6.55	6
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$K^{*+}K^{*0}; K^{*+} \longrightarrow K^0{}_S\pi^+; K^{*0} \longrightarrow K^-\pi^+$	3.57	4
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$\acute{\eta}\pi^+; \acute{\eta} ightarrow ho^0 \gamma$	4.26	
Total	50.79	



•Distribution of m_{BC} in data after unblinding in the $K^{*+}K^{*0}$; $K^{*+} \rightarrow K^0{}_S\pi^+$; $K^{*0} \rightarrow K^-\pi^+$ mode

Decay Mode of the D_S^+	Expected Yield from MC in 586 pb ⁻¹	Observed Yield in 586 pb ⁻¹
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$\pi^+\dot{\eta};\dot{\eta}{\rightarrow}\pi^+\pi^-\eta;\eta{\rightarrow}\gamma\gamma$	1.20	4
$K^{+}K^{-}\pi^{+}\pi^{0}$	6.55	6
$\pi^+\pi^-\pi^+$	5.32	7
$K^{*+}K^{*0}; K^{*+} \longrightarrow K^0{}_S\pi^+; K^{*0} \longrightarrow K^-\pi^+$	3.57	4
$\eta ho^+;\eta{ o}\gamma\gamma; ho^+{ o}\pi^+\pi^0$	8.11	7
$\acute{\eta}\pi^+$; $\acute{\eta}{ ightarrow} ho^0\gamma$	4.26	
Total	50.79	



Significance of Observation

Question: What is the probability that $12.61 \pm 1.58 \pm 1.53$ background events fluctuated up to the 51 events we observed?

Answer: 1.7×10⁻¹⁰

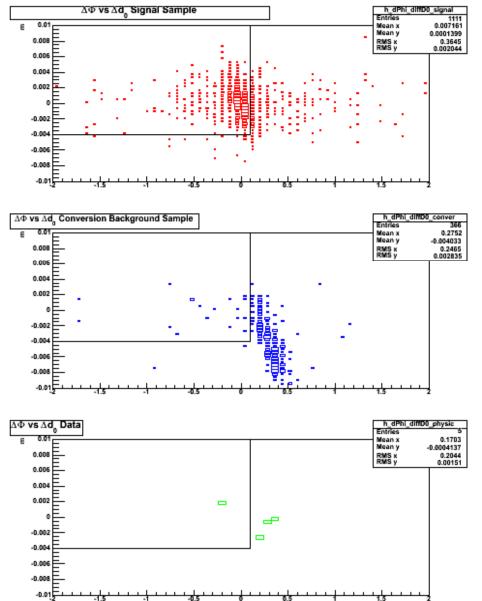
•Assume the estimated background fluctuates as a Gaussian with mean (*b*) = 12.61 and standard deviation (σ) = 2.20

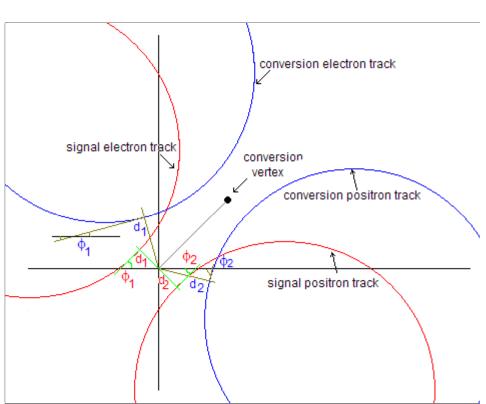
•Convolute it with a Poisson fluctuation to 51 events (*n*) or higher to find the probability.

$$P(b,\sigma,n) = \frac{\sum_{i=n}^{i=\infty} \int_{x=0}^{x=\infty} \frac{x^i}{i!} e^{-[x+\frac{1}{2}(\frac{x-b}{\sigma})^2]} dx}{\int_{x=0}^{x=\infty} e^{-\frac{1}{2}(\frac{x-b}{\sigma})^2} dx}$$

•In high energy physics, we express such small probabilities in terms of standard deviations of a Gaussian we need to eliminate to be left with that probability. 1.7×10^{-10} corresponds to 6.39 σ

$K^+K^-\pi^+$ Mode $\Delta\Phi$ vs Δd_0





The $\Delta \Phi \& \Delta d_0$ between the electron and positron in the signal (red) and conversion (blue)

Background Monte Carlo Samples

- The generic and continuum MC samples that come with the CLEO-c datasets do not have electron mass fitted tracks either!
- We do not bother reprocessing them. Instead we remove events in the generic MC that have $D_s^{*+} \rightarrow D_s^{+} \gamma$ at the level of event generation. We replace these events with privately produced $D_s^{*+} \rightarrow D_s^{+} \gamma$ where electron mass fitted tracks are stored.
- This complicates plots, hence not shown.
- However, such a separation of backgrounds is used for optimizing our selection criteria.

Criteria to Select $D_s^{*+} \rightarrow D_s^{+} \gamma$ Events

•Now we measure the yields and efficiencies for the $D_s^{*+} \rightarrow D_s^{+} \gamma$ channel since we have set out to measure the ratio of branching fractions

$$B(D_s^{*+} \to D_s^+ e^+ e^-) / B(D_s^{*+} \to D_s^+ \gamma)$$

•We reconstruct the D_s^{*+} through the D_s^{+} and the γ

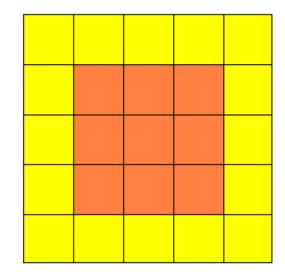
•The D_s^+ is reconstructed through the 9 hadronic decay modes

•Criteria on the e^+e^- inapplicable. Instead, criteria on the electromagnetic shower of the photon in the calorimeter are applied.

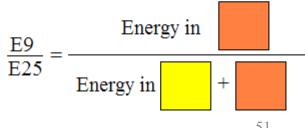
•10 MeV < Shower Energy < 2.0 GeV

- •No tracks leading to or in the vicinity of the shower
- •Known noisy calorimeter crystals discarded
- •EM shower shape ensured with *E9/E25*

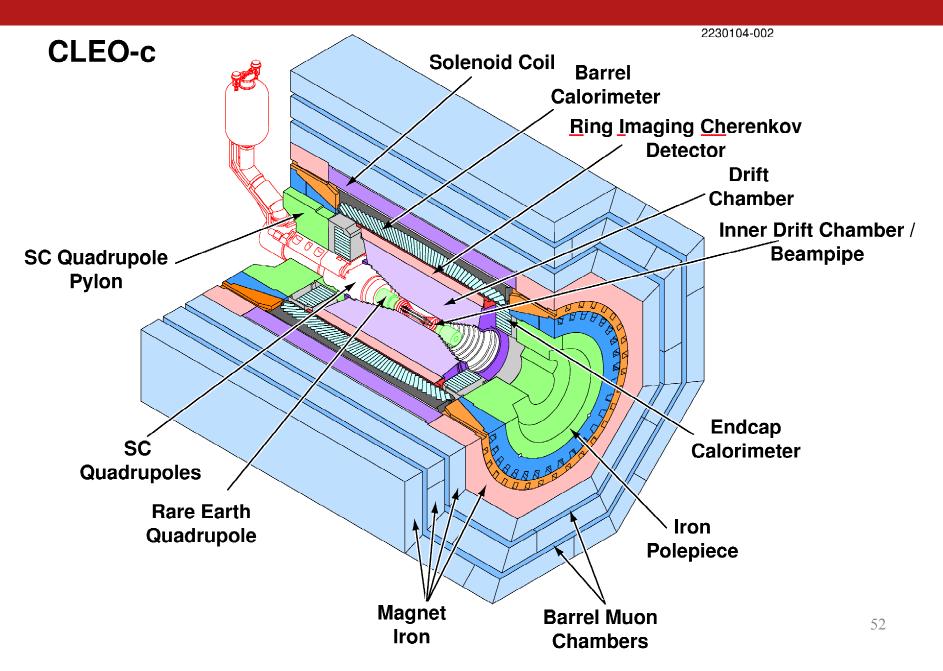
•The m_{BC} criterion is discarded in favor of a fit to extract yield.

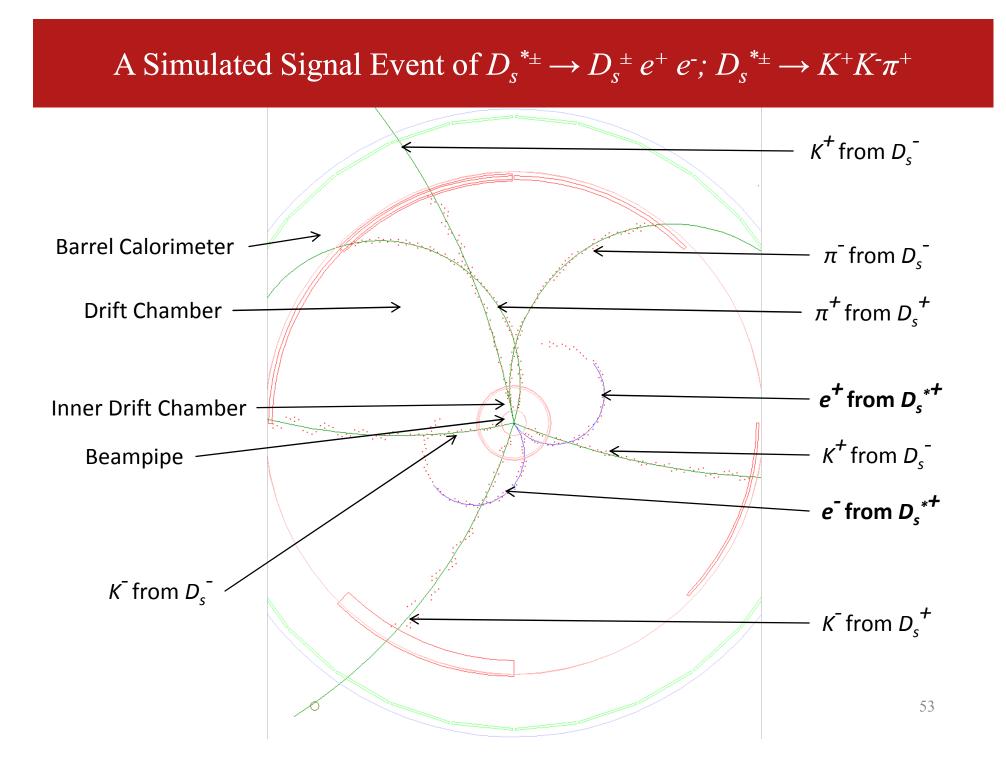


25 CsI crystals surrounding the center of an EM shower



The CLEO-c Detector





The Problem with Soft Electrons

- Hits in the CLEO-c drift chamber are fitted to a track with a Kalman filter
- Kalman filter depends on models of energy loss of charged particle in material
- Given a certain momentum, the energy loss depends on the mass of the particle
- Tracks associated with electrons in the standard CLEO-c dataset were Kalman fitted with the π^{\pm} mass. Fits with the e^{\pm} mass were computed but discarded in the interest of disk space.
- This works fine in reconstructing electrons above ~ 200 MeV, but we are working with electrons ~ 70 MeV each.
- The Kalman filter over-compensates the energy of electron tracks. Hurts our analysis.